

UWIS, Fluid Dynamics, Problem sheet 4

Thomas Kuster

11. Januar 2006

1 2D Flow field

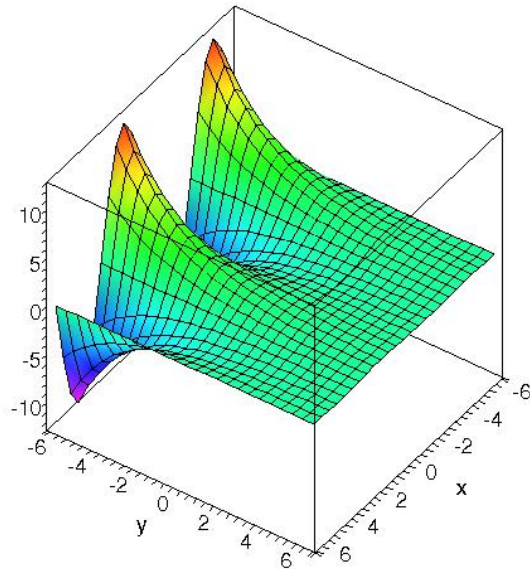
$$\psi = A \sin(kx) \exp(-ly)$$

1.1 Sketch

```
> restart;
> with(plots):

Warning, the name changecoords has been redefined
> psi := (x,y)->A * sin(k*x) * exp(-l*y);
      (x,y) ↦ sin(x) e-0.4y := (x,y) ↦ sin(x) e-0.4y
> A := 1;
k := 1;
l := 0.4;

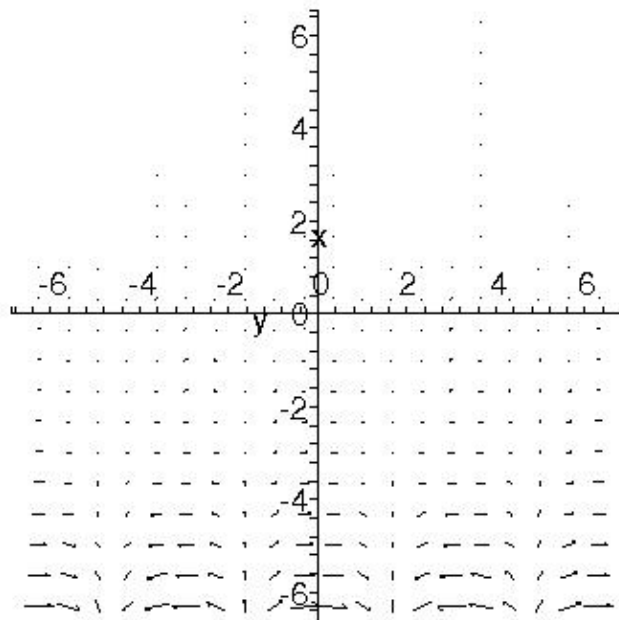
      A := 1
      k := 1
      l := 0.4
> plot3d(psi(x,y), x=-2*Pi..2*Pi, y=-2*Pi..2*Pi);
```



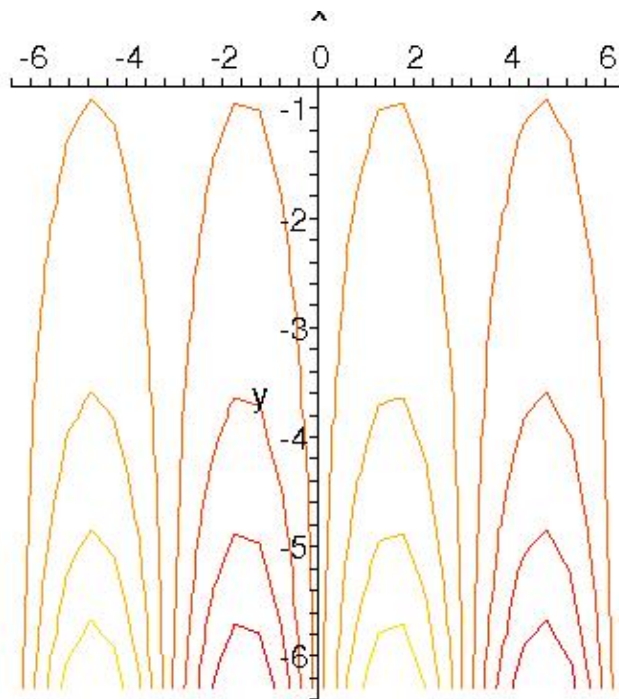
```
> with(plots):
```

```
Warning, the name changecoords has been redefined
```

```
> gradplot(psi(x,y), x=-2*Pi..2*Pi, y=-2*Pi..2*Pi);
```



```
> contourplot(psi(x,y), x=-2*Pi..2*Pi, y=-2*Pi..2*Pi);
```



1.2 Horizontal velocity

$$\frac{\partial \psi}{\partial y} = -Al \sin(kx) \exp(-ly) = u$$

$$\frac{\partial \psi}{\partial x} = Ak \cos(kx) \exp(-ly) = v$$

1.3 Vorticity

$$\zeta = \frac{\partial}{\partial x} v - \frac{\partial}{\partial y} u \tag{1}$$

$$\frac{\partial}{\partial x} v = -Ak^2 \sin(kx) \exp(-ly) \tag{2}$$

$$\frac{\partial}{\partial y} u = -Al^2 \sin(kx) \exp(-ly) \tag{3}$$

$$(1) \stackrel{(2)\&(3)}{\Rightarrow} \zeta = (l^2 - k^2) \underbrace{A \sin(kx) \exp(-ly)}_{\psi} \tag{4}$$

1.4 Streamfunction as a solution of a differential equation

$$\frac{D}{Dt}\zeta = \nu\Delta\zeta \quad (5)$$

$$\frac{D}{Dt}\zeta - \nu\Delta\zeta = 0 \quad |\zeta \stackrel{(4)}{=} (l^2 - k^2)\psi \quad (6)$$

$$\frac{D}{Dt}(l^2 - k^2)\psi - \nu\Delta((l^2 - k^2)\psi) = 0 \quad (7)$$

$$\underbrace{(l^2 - k^2)}_{=0 \Rightarrow l = \pm k} \left(\frac{D}{Dt}\psi - \nu\Delta\psi \right) = 0 \quad (8)$$

$$\frac{D}{Dt}\psi - \nu\Delta\psi = 0 \quad (9)$$

$$\vdots = \vdots \quad (10)$$

Now it is too late for good brain work, but here is the effort before Xmas.

Use ψ as basic approach for ζ :

$$\frac{\partial^2}{\partial x^2}u = -A \cos(kx)k^3 \exp(-ly)$$

$$\frac{\partial^2}{\partial y^2}u = A \cos(kx)kl^2 \exp(-ly)$$

$$\frac{\partial^2}{\partial x^2}v = A \sin(kx)k^2l \exp(-ly)$$

$$\frac{\partial^2}{\partial y^2}v = -A \sin(kx)l^3 \exp(-ly)$$

$$\Rightarrow \left(\begin{array}{c} \frac{\partial^2}{\partial x^2}u + \frac{\partial^2}{\partial y^2}u \\ \frac{\partial^2}{\partial x^2}v + \frac{\partial^2}{\partial y^2}v \end{array} \right) = \left(\begin{array}{c} -A \cos(kx)k^3 \exp(-ly) + A \cos(kx)kl^2 \exp(-ly) \\ A \sin(kx)k^2l \exp(-ly) - A \sin(kx)l^3 \exp(-ly) \end{array} \right) =$$

$$\left(\begin{array}{c} A \cos(kx) \exp(-ly)k(-k^2 + l^2) \\ A \sin(kx) \exp(-ly)l(k^2 - l^2) \end{array} \right)$$

$$\frac{D}{Dt} = 0 = \nu \left(\begin{array}{c} A \cos(kx) \exp(-ly)k(-k^2 + l^2) \\ A \sin(kx) \exp(-ly)l(k^2 - l^2) \end{array} \right)$$

Tree possibilites

$$\nu \left(\begin{array}{c} -k^2 + l^2 \\ k^2 - l^2 \end{array} \right) = 0 \Rightarrow k^2 - l^2 = 0 \Rightarrow k^2 = l^2 \Rightarrow k = \pm l$$

$$\left(\begin{array}{c} A \cos(kx) \exp(-ly)k(-k^2 + l^2) \\ A \sin(kx) \exp(-ly)l(k^2 - l^2) \end{array} \right) = 0$$

$$\begin{pmatrix} k \\ l \end{pmatrix} = 0 \quad (\text{trivial})$$

2 Flow through a tube: Law of Hagen Poiseuille

Navier-Stokes equation for the z-direction yields:

$$\frac{\partial}{\partial t}w + u \frac{\partial}{\partial r}w + \frac{v}{r} \frac{\partial}{\partial \varphi}w + w \frac{\partial}{\partial z}w = -\frac{1}{\rho} \frac{\partial}{\partial z}p + \nu \Delta w$$

2.1 Rearrange

$$\underbrace{\frac{\partial}{\partial t}w}_{=0 \text{ (steady state)}} + \underbrace{u}_{=0} \frac{\partial}{\partial r}w + \underbrace{\frac{v}{r}}_{=0} \frac{\partial}{\partial \varphi}w + w \underbrace{\frac{\partial}{\partial z}w}_{=0} = -\frac{1}{\rho} \frac{\partial}{\partial z}p + \nu \Delta w \quad (11)$$

$$0 = -\frac{1}{\rho} \frac{\partial}{\partial z}p + \nu \Delta w \quad (12)$$

$$\frac{1}{\rho} \frac{\partial}{\partial z}p = \nu \Delta w \quad (13)$$

$$\Delta w = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 w}{\partial \varphi^2} + \frac{\partial^2 w}{\partial z^2} \quad (14)$$

$$= \frac{1}{r} \left(\frac{\partial w}{\partial r} + r \frac{\partial^2 w}{\partial r^2} \right) + \underbrace{\frac{1}{r^2} \frac{\partial^2 w}{\partial \varphi^2}}_{=0 \text{ axial-symmetric}} + \underbrace{\frac{\partial^2 w}{\partial z^2}}_{=0 \text{ no source}} \quad (15)$$

$$(13) \stackrel{(15)}{\Rightarrow} \frac{1}{\rho} \frac{\partial p}{\partial z} = \nu \frac{1}{r} \left(\frac{\partial w}{\partial r} + r \frac{\partial^2 w}{\partial r^2} \right) \quad (16)$$

2.2 Velocity profile

$$\frac{\partial p}{\partial z} = \text{const} = c \quad (17)$$

$$(16) \Rightarrow \underbrace{\frac{1}{\rho} \frac{\partial p}{\partial z}}_{=c} = \nu \frac{1}{r} \frac{\partial w}{\partial r} + \nu \frac{\partial^2 w}{\partial r^2} \quad \Bigg| \cdot \frac{r}{\nu} \quad (18)$$

$$\frac{cr}{\rho\nu} = \frac{\partial w}{\partial r} + r \frac{\partial^2 w}{\partial r^2} \quad \Bigg| \text{“from (15) to (14)”} \quad (19)$$

$$\frac{cr}{\rho\nu} = \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) \quad \Bigg| \underbrace{\int_0^r}_{\text{integrate over r from 0 to r, that is unclear}} dr \quad (20)$$

$$\int_0^r \frac{cr}{\rho\nu} dr = \int_0^r \frac{\partial}{\partial r} \left(r \frac{\partial w}{\partial r} \right) dr \quad (21)$$

$$\frac{cr^2}{2\rho\nu} \Bigg|_0^r = r \frac{\partial w}{\partial r} \Bigg|_0^r \quad (22)$$

$$\frac{cr^2}{2\rho\nu} = r \frac{\partial w}{\partial r} \quad \Bigg| \cdot \frac{\partial r}{r} \quad (23)$$

$$\frac{cr}{2\rho\nu} \partial r = \partial w \quad \Bigg| \int \quad (24)$$

$$\int_0^r \frac{cr}{2\rho\nu} \partial r = \int_{w(0)}^{w(r)} \partial w \quad (25)$$

$$\frac{cr^2}{4\rho\nu} = w(r) - w(0) \quad (26)$$

$$\frac{cr^2}{4\rho\nu} + w(0) = w(r) \quad \Bigg| \text{boundary condition } w(R) = 0 \quad (27)$$

$$\frac{cR^2}{4\rho\nu} + w(0) = 0 \quad (28)$$

$$w(0) = -\frac{cR^2}{4\rho\nu} \quad (29)$$

$$(27) \stackrel{(29)}{\Rightarrow} w(r) = \frac{cr^2}{4\rho\nu} - \frac{cR^2}{4\rho\nu} \quad (30)$$

$$w(r) = \frac{c}{4\rho\nu} (r^2 - R^2) \quad (31)$$

$$w(r) \stackrel{(17)}{=} \frac{1}{4\rho\nu} \frac{\partial p}{\partial z} (r^2 - R^2) \quad (32)$$

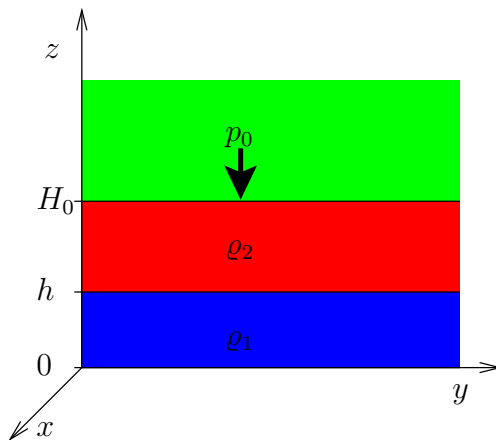
2.3 Mean flux

$$\begin{aligned}
 & \int_0^{2\pi} \int_0^R w(r) r dr d\varphi = \\
 & \frac{1}{4\rho\nu} \frac{\partial p}{\partial z} \int_0^{2\pi} \int_0^R (r^2 - R^2) r dr d\varphi = \\
 & \frac{1}{4\rho\nu} \frac{\partial p}{\partial z} \int_0^{2\pi} \int_0^R r^3 - R^2 r dr d\varphi = \\
 & \frac{1}{4\rho\nu} \frac{\partial p}{\partial z} \int_0^{2\pi} \left(\frac{r^4}{4} - \frac{R^2 r^2}{2} \right) \Big|_0^R d\varphi = \\
 & \frac{1}{4\rho\nu} \frac{\partial p}{\partial z} \int_0^{2\pi} \left(\frac{R^4}{4} - \frac{R^4}{2} \right) d\varphi = \\
 & \frac{1}{4\rho\nu} \frac{\partial p}{\partial z} \int_0^{2\pi} -\frac{R^4}{4} d\varphi = \\
 & \frac{1}{4\rho\nu} \frac{\partial p}{\partial z} \left(-\frac{R^4}{4} \varphi \right) \Big|_0^{2\pi} = \\
 & -\frac{1}{4\rho\nu} \frac{\partial p}{\partial z} \frac{R^4}{4} 2\pi = \\
 & \quad -\frac{R^4 \pi}{8\rho\nu} \frac{\partial p}{\partial z}
 \end{aligned}$$

2.4 Derive of the mean velocity

$$\begin{aligned}
 \bar{w} &= \frac{F}{A} \\
 &= \frac{-\frac{R^4 \pi}{8\rho\nu} \frac{\partial p}{\partial z}}{R^2 \pi} \\
 &= -\frac{R^2}{8\rho\nu} \frac{\partial p}{\partial z}
 \end{aligned}$$

3 Shallow water system



3.1 Pressure gradient force

$$\begin{aligned} \frac{dp}{dz} &= -\rho g \\ dp &= -\rho g dz \\ \int_{p(0)}^{p(h)} dp &= \int_0^h -\rho_1 g dz \\ p(h) - p(0) &= -\rho_1 g(h - 0) \\ -p(0) &= -\rho_1 gh - \underbrace{p(h)}_{(H_0-h)\rho_2 g + p_0} \\ p(0) &= \rho_1 gh + (H_0 - h)\rho_2 g + p_0 \\ p(0) &= \rho_1 gh - \rho_2 gh + \rho_2 g H_0 + p_0 \\ p(0) &= gh(\rho_1 - \rho_2) + \underbrace{\rho_2 g H_0 + p_0}_{=0} \\ \frac{p(0)}{\rho_1} &= gh(\rho_1 - \rho_2) \\ \frac{p(0)}{\rho_1} &= gh \frac{\rho_1 - \rho_2}{\rho_1} \\ \frac{p(0)}{\rho_1} &\stackrel{g^* = g \frac{\rho_1 - \rho_2}{\rho_1}}{=} g^* h \\ \frac{1}{\rho_1} p_x &= g^* h_x \end{aligned}$$

3.2 Non-dimensional

$$u_t + uu_x + vu_x - fv = -g^* h_x \quad (33)$$

$$v_t + uv_x + vv_y + fu = -g^* h_y \quad (34)$$

$$h_t + uh_x + vh_y + hu_x + hv_y = 0 \quad (35)$$

$$\begin{aligned}
 u &= \tilde{u}U \\
 u_x &= \tilde{u}_x \frac{U}{L} \\
 u_y &= \tilde{u}_y \frac{U}{L} \\
 u_t &= \tilde{u}_t \frac{U^2}{L} \\
 v &= \tilde{v}V = \tilde{v}U \\
 v_x &= \tilde{v}_x \frac{V}{L} = \tilde{v}_x \frac{U}{L} \\
 v_y &= \tilde{v}_y \frac{V}{L} = \tilde{v}_y \frac{U}{L} \\
 h &= \tilde{h}H \\
 h_x &= \tilde{h}_x \frac{H}{L} \\
 h_y &= \tilde{h}_y \frac{H}{L} \\
 h_t &= \tilde{h}_t \frac{HU}{L} \\
 R_0 &= \frac{U}{fL} \Rightarrow R_0 \frac{fL}{U} = 1
 \end{aligned}$$

$$\begin{aligned}
 (33) \Rightarrow u_t + uu_x + vu_y - fv &= -g^* h_x \\
 \tilde{u}_t \frac{U^2}{L} + \tilde{u}U \tilde{u}_x \frac{U}{L} + \tilde{v}U \tilde{u}_y \frac{U}{L} - f\tilde{v}U &= -g^* \tilde{h}_x \frac{H}{L} \\
 R_0 \frac{fL}{U} \left(\tilde{u}_t \frac{U^2}{L} + \tilde{u}U \tilde{u}_x \frac{U}{L} + \tilde{v}U \tilde{u}_y \frac{U}{L} \right) - R_0 \frac{fL}{U} f\tilde{v}U &= -g^* \tilde{h}_x \frac{H}{L} R_0 \frac{fL}{U} \quad \Big| : f \\
 R_0 (\tilde{u}_t U + \tilde{u}U \tilde{u}_x + \tilde{v}U \tilde{u}_y) - R_0 L f \tilde{v} &= -g^* \tilde{h}_x R_0 \frac{H}{U} \quad \Big| R_0 = \frac{U}{fL} \\
 R_0 U (\tilde{u}_t + \tilde{u} \tilde{u}_x + \tilde{v} \tilde{u}_y) - U \tilde{v} &= -g^* \tilde{h}_x R_0 \frac{H}{U} \quad \Big| : U \\
 R_0 (\tilde{u}_t + \tilde{u} \tilde{u}_x + \tilde{v} \tilde{u}_y) - \tilde{v} &= \underbrace{-g^* \frac{H}{U^2}}_{=\exists} R_0 \tilde{h}_x
 \end{aligned}$$

$$\begin{aligned}
 (34) \Rightarrow v_t + uv_x + vv_y + fu &= -g^* h_y \\
 R_0 \frac{fL}{U} \left(\tilde{v}_t \frac{U^2}{L} + \tilde{u}U\tilde{v}_x \frac{U}{L} + \tilde{v}U\tilde{v}_x \frac{U}{L} \right) + R_0 \frac{fL}{U} f\tilde{u}U &= -g^* \tilde{h}_y \frac{H}{L} R_0 \frac{fL}{U} \quad | : f \\
 R_0 (\tilde{v}_t U + \tilde{u}U\tilde{v}_x + \tilde{v}U\tilde{v}_y) + R_0 L f \tilde{u} &= -g^* \tilde{h}_y R_0 \frac{H}{U} \quad | R_0 = \frac{U}{fL} \\
 R_0 U (\tilde{v}_t + \tilde{u}\tilde{v}_x + \tilde{v}\tilde{v}_y) + U \tilde{u} &= -g^* \tilde{h}_y R_0 \frac{H}{U} \quad | : U \\
 R_0 (\tilde{v}_t + \tilde{u}\tilde{v}_x + \tilde{v}\tilde{v}_y) + \tilde{u} &= \underbrace{-g^* \frac{H}{U^2} R_0 \tilde{h}_y}_{=\exists}
 \end{aligned}$$

$$\begin{aligned}
 (35) \Rightarrow h_t + uh_x + vh_y + hu_x + hv_y &= 0 \\
 \tilde{h}_t \frac{HU}{L} + \tilde{u}U\tilde{h}_x \frac{H}{L} + \tilde{v}U\tilde{h}_y \frac{H}{L} + \tilde{h}H\tilde{u}_x \frac{U}{L} + \tilde{h}H\tilde{v}_y \frac{U}{L} &= 0 \quad | : \frac{HU}{L} \\
 \tilde{h}_t + \tilde{u}\tilde{h}_x + \tilde{v}\tilde{h}_y + \tilde{h}\tilde{u}_x + \tilde{h}\tilde{v}_y &= 0
 \end{aligned}$$

3.3 Dimensionless parameter

Why you use “ \exists ” as a parameter? For me is “ \exists ” reserved for “it exists”.
 From (3.2):

$$\exists = -g^* \frac{H}{U^2}$$

3.4 R_0 oft the observed motions

$$R_0 = \frac{U}{fL}$$

High and Lows	$R_0 = \frac{10}{10^{-4} \cdot 10^6} = 0.1$	Coriolis effect plays the major rule
Hurricanes	$R_0 = \frac{10}{10^{-4} \cdot 10^5} = 1$	$\frac{du}{dt}$ and coriolis effect are equated
Tornados	$R_0 = \frac{10^2}{10^{-4} \cdot 10^3} = 1000$	$\frac{du}{dt}$ plays the major rule

3.5 \exists for high and lows

$$\exists = -g^* \frac{H}{U^2} = -g^* \frac{H}{U^2} = -10 \frac{\text{m}}{\text{s}^2} \frac{10^2 \text{m}}{10^2 \frac{\text{m}^2}{\text{s}^2}} = -10$$